

Kochen-Specker theorem and games

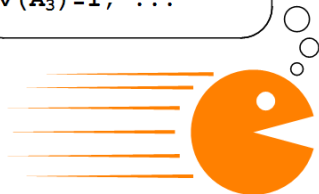
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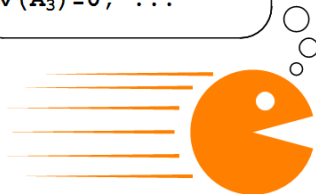
December 13, 2007

Hidden variables

$$\begin{aligned}v(A_1) &= 0, & v(A_2) &= 1, \\v(A_3) &= 1, & \dots\end{aligned}$$



$$\begin{aligned}v(A_1) &= 1, & v(A_2) &= 1, \\v(A_3) &= 0, & \dots\end{aligned}$$



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In this talk

We will consider **proofs** of several versions of Kochen-Specker theorem and **games** that are based on these proofs.

Observables

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Measurement described by an observable

Observable M is a **Hermitian** operator. If

$$M = \sum \lambda P_\lambda$$

is a spectral decomposition of M , then M defines a projective measurement in the following way:

- the outcome of the measurement is an **eigenvalue** λ of M ,
- the state collapses to the corresponding **eigenspace** P_λ .

Commuting observables

Definition

Observables A and B are said to **commute** if

$$AB = BA$$

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Theorem

If mutually commuting observables A_1, A_2, \dots, A_n satisfy some functional identity

$$f(A_1, A_2, \dots, A_n) = 0,$$

then the values assigned to them in an individual system must also be related by

$$f\left(v(A_1), v(A_2), \dots, v(A_n)\right) = 0$$

Kochen-Specker theorem (3 dimensional version)

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Consequences of Kochen-Specker theorem

Every non-contextual hidden variables theory is inconsistent with quantum mechanics formalism.

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- Observable S_v measures the square of spin component of a spin 1 particle along direction $v \in \mathbb{R}^3$

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- The outcomes (eigenvalues) of the measurement S_v are 1 or 0
- If $\{u, v, w\}$ are mutually orthogonal vectors in \mathbb{R}^3 , then
 - ① $\{S_u, S_v, S_w\}$ is a set of mutually commuting observables
 - ② $S_u + S_v + S_w = 2I$

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 - ① $\{S_u, S_v, S_w\}$ is a set of mutually commuting observables
 - ② $S_u + S_v + S_w = 2I \implies v(S_u) + v(S_v) + v(S_w) = 2.$

The task of proving Kochen-Specker theorem can be reduced to the following problem

Find a set of vectors in \mathbb{R}^3 for which it is impossible to assign “0” and “1” (outcomes of observables S_v) so that in each set of three mutually orthogonal vectors “1” is assigned to exactly two of them.

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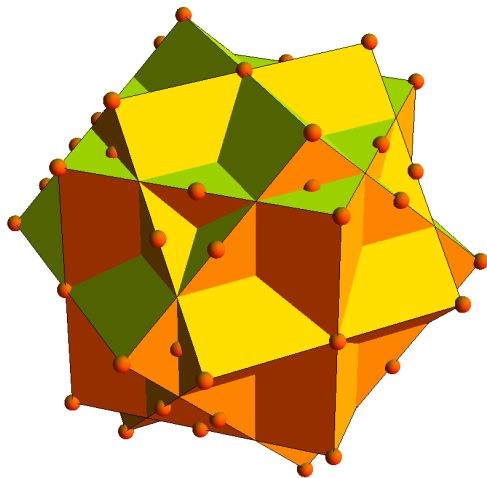
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- 1 Kochen and Specker (1967) found the required set with **117** vectors
- 2 Later Conway and Kochen reduced the set to **31** vectors
- 3 Peres (1991) found the required set with **33** vectors (with nice **symmetries**)

Magic configuration



Although it is not obvious, this set satisfies the required property.

M.C. Escher "Waterfall"



Kochen-Specker game

Setting of the game

- Alice and Bob plays against verifier

Kochen-Specker game

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As always we will see that **entanglement** turns out to be the key trick in quantum strategy.

Kochen-Specker game

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- Verifier chooses three mutually orthogonal vectors v_i, v_j, v_k from the set V . He asks
 - Alice to assign “0” or “1” to each of these vectors
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 - **Parity rule:** “1” gets assigned to exactly two of the three vectors
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 - **Consistency rule**: Alice and Bob assigns the same values to vector v_l
- Alice and Bob cannot always win if they use **classical** strategy as this would lead to violation of KS theorem.
- Yet they can win using **quantum** strategy with **entanglement**.

Quantum strategy for KS game

Alice and Bob share the state $|\Psi\rangle = \frac{1}{\sqrt{3}}(|00\rangle + |11\rangle + |22\rangle)$.

- 1 Alice measures her qutrit with POVM $\{|v_i\rangle\langle v_i|, |v_j\rangle\langle v_j|, |v_k\rangle\langle v_k|\}$. She assigns “0” to the vector corresponding to the outcome of her measurement and “1” to the rest two vectors.
- 2 Bob measures with POVM $\{|v_l\rangle\langle v_l|, I - |v_l\rangle\langle v_l|\}$. He assigns “0” to vector v_l if the state collapses to $|v_l\rangle$ and “1” if otherwise.

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We need to check whether **parity** and **consistency** rules are always satisfied.

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- 4 dimensional Hilbert space corresponds to **two qubit** system
- Again we will construct a set of observables that satisfy some functional identities that cannot be satisfied by the values assigned to them.

Magic square

	X	Y	Z
X	I	iZ	$-iY$
Y	$-iZ$	I	iX
Z	iY	$-iX$	I

	$-I$	$-I$	$-I$	
$I \otimes Z$	$Z \otimes I$	$Z \otimes Z$		I
$X \otimes I$	$I \otimes X$	$X \otimes X$		I
$-X \otimes Z$	$-Z \otimes X$	$Y \otimes Y$		I

Multiplication of Pauli matrices.

Magic square.

- Observables on each row and column are **mutually commuting**.

Magic square

	X	Y	Z
X	I	iZ	$-iY$
Y	$-iZ$	I	iX
Z	iY	$-iX$	I

	-I	-I	-I	
I	$I \otimes Z$	$Z \otimes I$	$Z \otimes Z$	I
X	$X \otimes I$	$I \otimes X$	$X \otimes X$	I
Y	$-X \otimes Z$	$-Z \otimes X$	$Y \otimes Y$	I

Multiplication of Pauli matrices.

Magic square.

- Observables on each row and column are **mutually commuting**.
- It is **impossible** to fill in the outcomes of observables so that functional identities are satisfied.

Game

- Verifier asks Alice to fill in some row and Bob to fill some column with “1” or “-1”

	$-I$	$-I$	$-I$	
$I \otimes Z$	$Z \otimes I$	$Z \otimes Z$		I
$X \otimes I$	$I \otimes X$	$X \otimes X$		I
$-X \otimes Z$	$-Z \otimes X$	$Y \otimes Y$		I

Game

- Verifier asks Alice to fill in some row and Bob to fill some column with “1” or “-1”
- Alice and Bob win if
 - **Parity rule** The parity of “-1” is even for Alice and odd for Bob
 - **Consistency rule** Alice and Bob assign the same value to the intersection

	-I	-I	-I	
I \otimes Z	Z \otimes I	Z \otimes Z		I
X \otimes I	I \otimes X	X \otimes X		I
-X \otimes Z	-Z \otimes X	Y \otimes Y		I

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- There is no **perfect classical** strategy.

	-I	-I	-I	
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	-I	-I	-I	
I⊗Z	Z⊗I	Z⊗Z		I
X⊗I	I⊗X	X⊗X		I
-X⊗Z	-Z⊗X	Y⊗Y		I

Quantum strategy

- Alice and Bob share $|\Psi\rangle = \left(\frac{1}{\sqrt{2}}(|00\rangle + |11\rangle)\right)^{\otimes 2}$

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Quantum strategy

- Alice and Bob share $|\Psi\rangle = \left(\frac{1}{\sqrt{2}}(|00\rangle + |11\rangle)\right)^{\otimes 2}$
- Alice (Bob) measures her part of $|\Psi\rangle$ with the observables on the corresponding row (column) and gives verifier the outcomes of her measurement.

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	-1	-1	-1	
I⊗Z	Z⊗I	Z⊗Z		I
X⊗I	I⊗X	X⊗X		I
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We have to check whether **parity** and **consistency** rules hold.

Consistency rule verification

- Let $\mathcal{B} = \{|b_1\rangle, |b_2\rangle, |b_3\rangle, |b_4\rangle\}$ be a basis of Alice's state space (2 qubits) and $\overline{\mathcal{B}} = \{\overline{|b_1\rangle}, \overline{|b_2\rangle}, \overline{|b_3\rangle}, \overline{|b_4\rangle}\}$ be a basis of Bob's state space.

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- It turns out that $|\Psi\rangle = \left(\frac{1}{\sqrt{2}}(|00\rangle + |11\rangle)\right)^{\otimes 2}$ in these basis can be written as:

$$|\Psi\rangle = \frac{1}{4} \left(|b_1\rangle |\overline{b_1}\rangle + |b_2\rangle |\overline{b_2}\rangle + |b_3\rangle |\overline{b_3}\rangle + |b_4\rangle |\overline{b_4}\rangle \right)$$

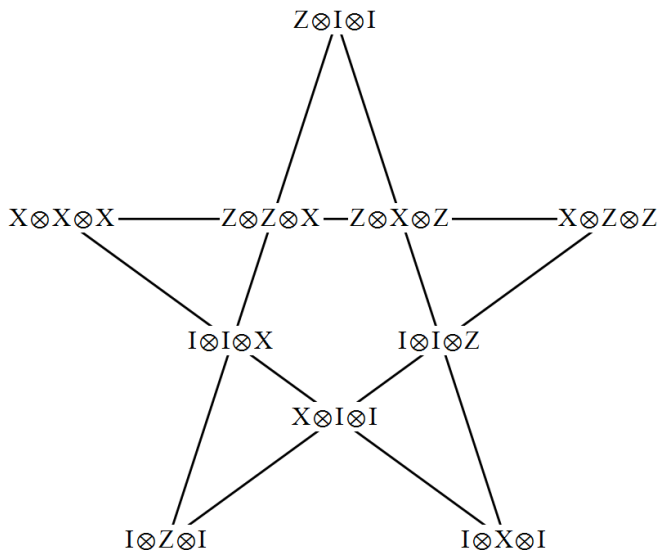
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- Also it can be shown that the eigenvectors of observables being measured are real, therefore $\mathcal{B} = \overline{\mathcal{B}}$ and Bob will get the same outcome as Alice.

Magic star





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